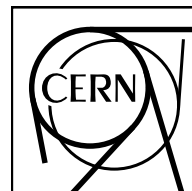


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## Search for microscopic black holes in a like-sign dimuon final state using large track multiplicity with the ATLAS detector

The ATLAS Collaboration

### Abstract

A search is presented for microscopic black holes in a like-sign dimuon final state in proton–proton collisions at  $\sqrt{s} = 8$  TeV. The data were collected with the ATLAS detector at the Large Hadron Collider in 2012 and correspond to an integrated luminosity of  $20.3 \text{ fb}^{-1}$ . Using a high track multiplicity requirement,  $0.6 \pm 0.2$  background events from Standard Model processes are predicted and none observed. This result is interpreted in the context of low-scale gravity models and 95% CL lower limits on microscopic black hole masses are set for different model assumptions.

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The ATLAS Collaboration  
(Dated: August 20, 2013)

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## I. INTRODUCTION

The hierarchy problem, in which the Planck scale ( $M_{\text{Pl}} \approx 10^{19} \text{ GeV}$ ) is much higher than the electroweak scale ( $\approx 100 \text{ GeV}$ ), provides a strong motivation to search for new phenomena not described by the Standard Model of particle physics. A theory introducing extra dimensions is one possible solution. In some models of extra dimensions, the gravitational field propagates into  $n + 4$  dimensions, where  $n$  is the number of extra dimensions beyond the four space-time dimensions. One of the models of extra dimensions is the model proposed by Arkani-Hamed, Dimopoulos, and Dvali (ADD) [1–3], in which the gravitational field propagates into large, flat, extra dimensions while the Standard Model particles are localized in four space-time dimensions. Since the gravitational field propagates into the extra dimensions, it is measured at a reduced strength in the four space-time dimensions. Thus, the fundamental Planck scale in  $D = 4 + n$  dimensions,  $M_D$ , could be comparable with the electroweak scale.

If extra dimensions exist and  $M_D$  is of the order of 1 TeV, microscopic black holes with TeV-scale mass could exist and be produced at the Large Hadron Collider (LHC) [4–8]. These black holes are produced when the impact parameter of the two colliding protons is smaller than the higher-dimensional event horizon of a black hole with mass equal to the invariant mass of the colliding proton system.

The black hole production has a continuous mass distribution ranging from  $M_D$  to the proton–proton center-of-mass energy. The black holes evaporate by emitting Hawking radiation [9], which determines the energy and multiplicity of the emitted particles. The relative multiplicities of different particle types are determined by the number of degrees of freedom of each particle type and the decay modes of the emitted unstable particles. Black hole events are thus expected to have a high multiplicity of high-momentum particles.

This paper describes a search for black holes in a like-sign dimuon final state. This final state can arise from muons directly produced by the black hole, or from the decay of Standard Model particles produced by the black

hole. The final state is expected to have low Standard Model backgrounds while retaining a high signal acceptance. Since the microscopic black holes can decay to a large number of particles with high transverse momentum ( $p_T$ ), the total track multiplicity of the event is exploited to distinguish signal events from backgrounds. The final result is obtained from the event yield in a signal region defined by high track multiplicity.

The following assumptions and conventions apply in this analysis. The classical approximations used for black hole production, and the semi-classical approximations for the decay are predicted to be valid only for black hole masses well above  $M_D$ . A lower threshold ( $M_{\text{TH}}$ ) is applied to the black hole mass,  $M_{\text{TH}} > M_D + 0.5 \text{ TeV}$ , to reduce contributions from regions where the models are invalid, and the production cross section is set to zero if the proton–proton center of mass energy is below  $M_{\text{TH}}$ . The mass of the produced black hole decreases from  $M_{\text{TH}}$  to  $M_D$  as a result of the emission of Hawking radiation. When the mass of the black hole approaches  $M_D$ , quantum gravity effects become important. In the final stage of the black hole decay, classical evaporation is no longer a good description. In such cases where the black hole mass is near  $M_D$ , the burst model adopted by the BLACKMAX event generator [10] is used in the final part of the decay. No graviton initial-state radiation or emission from the black hole is considered in this paper. Models of rotating and non-rotating black holes are both studied. The track multiplicity is predicted to be slightly lower for rotating black holes [11].

A previous result from ATLAS [12] in this final state excludes at 95% confidence level (CL) the production of black holes with  $M_{\text{TH}} \leq 3.3, 3.6$  and  $3.7 \text{ TeV}$  for  $M_D$  of 1.5 TeV and for  $n = 2, 4$ , and 6, respectively. A previous search by the ATLAS Collaboration in a lepton+jets final state [13] excludes at 95% CL black holes with  $M_{\text{TH}} \leq 4.5 \text{ TeV}$  for  $M_D = 1.5 \text{ TeV}$  and  $n = 6$ . The CMS Collaboration has conducted searches in a multi-object final state and excluded the production of black holes at 95% CL with  $M_{\text{TH}} \leq 5.7, 6.1$ , and  $6.2 \text{ TeV}$  for  $M_D = 1.5 \text{ TeV}$  and for  $n = 2, 4$ , and 6, respectively [14, 15].

The rest of this paper is organized as follows. A brief

description of the ATLAS detector is given in Sec. II. The data and simulation samples are described in Sec. III, followed by the event selection in Sec. IV. The background estimation techniques are discussed in Sec. V. The final results and their interpretation are presented in Sec. VI.

## II. THE ATLAS DETECTOR

The ATLAS detector [16] is a multi-purpose detector with a forward-backward symmetric cylindrical geometry covering nearly the entire solid angle [49] around the collision point with layers of tracking detectors, calorimeters and muon chambers. The inner detector is immersed in a 2 T axial magnetic field, provided by a solenoid, in the  $z$  direction and provides charged particle tracking in the pseudorapidity range  $|\eta| < 2.5$ . A silicon pixel detector covers the luminous region and typically provides three measurements per track, followed by a silicon microstrip tracker (SCT) that provides measurements from eight strip layers. In the region with  $|\eta| < 2.0$ , the silicon detectors are complemented by a transition radiation tracker (TRT), which provides more than 30 straw-tube measurements per track.

The calorimeter system covers the range  $|\eta| < 4.9$ . Lead/liquid argon (lead/LAr) electromagnetic sampling calorimeters cover the range  $|\eta| < 3.2$ , with an additional thin lead/LAr presampler covering  $|\eta| < 1.8$  to correct for energy loss in material upstream of the calorimeters. Hadronic calorimetry is provided by a steel/scintillator-tile calorimeter over  $|\eta| < 1.7$  and two copper/LAr endcap calorimeters over  $1.75 < |\eta| < 3.2$ . The solid angle coverage is completed with forward copper/LAr and tungsten/LAr calorimeters for electromagnetic and hadronic measurements respectively up to  $|\eta|$  of 4.9.

The muon spectrometer consists of separate trigger and high-precision tracking chambers that measure the deflection of muon tracks in a magnetic field with a bending integral in the range of 2 Tm to 8 Tm. The magnetic field is generated by three superconducting air-core toroid magnet systems. The tracking chambers cover the region  $|\eta| < 2.7$  with three layers of monitored drift tubes supplemented by cathode strip chambers in the innermost region of the endcap muon spectrometer. The muon trigger system covers the range  $|\eta| < 2.4$  with resistive plate chambers in the barrel, and thin gap chambers in the endcap regions.

## III. DATA AND MONTE CARLO SAMPLES

The data used in this analysis were collected with the ATLAS detector from proton-proton collisions produced at  $\sqrt{s} = 8$  TeV in 2012. The data correspond to an integrated luminosity of  $20.3 \text{ fb}^{-1}$ . The uncertainty on the luminosity is 2.8% and is derived, following the same methodology as that detailed in Ref. [17], from a preliminary calibration of the luminosity scale ob-

tained from beam-separation scans performed in November 2012. The events used for this analysis were recorded with a single-muon trigger with a threshold at 36 GeV on the muon  $p_T$ . The single muon trigger efficiency reaches a plateau for muons with  $p_T > 40$  GeV and the plateau efficiency is 71% in the barrel and 87% in the endcap for muons reconstructed offline. The inefficiency in the trigger is driven mainly by the uninstrumented regions of the muon trigger system.

Monte Carlo (MC) samples are used for both signal and background modeling. The ATLAS detector is simulated using GEANT4 [18], and simulation samples [19] are reconstructed using the same software as that used for the collision data. The effect of additional proton-proton collisions in the same or neighboring bunch crossings is modeled by overlaying simulated minimum-bias events onto the original hard-scattering event. MC events are then re-weighted so that the reconstructed vertex multiplicity distribution agrees with the one from data.

The dominant background processes are top-quark pair ( $t\bar{t}$ ), diboson, and  $W$ +jets production with smaller contributions from single-top production. Background MC samples for the  $t\bar{t}$  and the single-top ( $Wt$ -channel) processes are generated using POWHEG [20] and the CT10 [21] parton distribution functions (PDFs). Fragmentation and hadronization of the events is done with PYTHIA v6.426 [22] using the Perugia tune [23]. The top-quark mass is fixed at 172.5 GeV. Alternative samples for studying the systematic uncertainty are made using the ALPGEN v2.14 [24] or MC@NLO v4.03 [25] generators with HERWIG v6.520 [26] used for hadronization and JIMMY v4.31 [27] used to model the underlying event for both generators. The nominal single-top sample uses the diagram-removal scheme [28] and an alternative sample using the diagram-subtraction scheme is produced for systematic studies. Diboson samples ( $WZ$  and  $ZZ$ ) are generated and hadronized using SHERPA v1.4.1 [29]. The diboson samples are produced with the CT10 PDF set and use the ATLAS Underlying Event Tune 2B (AUET2B) [30]. These samples include the case where the  $Z$  boson (or  $\gamma^*$ ) is off-shell, with the invariant mass of the  $\gamma^*$  required to be above twice the muon mass.

Signal MC samples are generated using BLACKMAX v2.02 [10, 31] and the MSTW 2008 LO [32] PDF set with the mass of the black hole used as the factorization and renormalization scale. The signal samples are hadronized with PYTHIA v8.165 [33] using the AUET2B tune. Signal samples for rotating and non-rotating black holes are produced by varying  $M_D$  between 1.0 TeV and 4.5 TeV, and  $M_{TH}$  between 3.0 TeV and 6.5 TeV. In each case, samples are generated with  $n = 2, 4$ , and 6. As an illustration, the expected yield from rotating black holes in a model with  $n = 4$ ,  $M_{TH} = 5$  TeV, and  $M_D = 1.5$  TeV is shown throughout the paper.

#### IV. EVENT SELECTION

Events in data passing the single-muon trigger are selected for this analysis. The detector was required to have been operating properly when these events were collected. Events are also required to have a primary vertex reconstructed from at least five tracks with  $p_T > 400$  MeV. In events with multiple vertices, the vertex whose associated tracks have the largest  $\Sigma p_T^2$  is identified as the primary vertex.

Muon candidates are reconstructed from tracks measured in the muon spectrometer (MS). The MS tracks are matched with inner detector (ID) tracks using a procedure that takes material effects into account. The final parameters for the muon candidates are obtained from a statistical combination of the measured quantities in the MS and the ID. The muon candidates must satisfy  $|\eta| < 2.4$  and have  $p_T > 15$  GeV. The quality of the ID track associated with a muon is ensured by imposing requirements [34] on the number of pixel, SCT and TRT hits associated with the track. The ID tracks must also pass a requirement on the longitudinal impact parameter ( $z_0$ ) with respect to the primary vertex,  $|z_0 \sin \theta| < 1.5$  mm.

Events are required to have at least two muons. The two muons with the highest  $p_T$  are required to have the same charge. The muon with the highest  $p_T$  in the event is called the leading muon, while the muon with the second highest  $p_T$  is called the subleading muon. The leading muon is required to satisfy  $p_T > 40$  GeV to be above the trigger threshold and pass requirements on isolation and transverse impact parameter as described below. No such requirements are made for the subleading muon.

The muon isolation is constructed from the sum of transverse momenta of other ID tracks in a cone in  $\eta$ - $\phi$  space of radius  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.2$  around the muon. For the leading muon, the sum is required to be less than 20% of the muon  $p_T$ . The impact parameter significance for the muons is defined as  $|d_0/\sigma(d_0)|$ , where  $d_0$  is the transverse impact parameter of the muon, and  $\sigma(d_0)$  is the associated uncertainty. The leading muon must satisfy  $|d_0/\sigma(d_0)| \leq 3.0$ . The leading and subleading muons are required to be separated by  $\Delta R > 0.2$ .

The total track multiplicity ( $N_{\text{trk}}$ ) of the event is calculated by considering all ID tracks with  $p_T > 10$  GeV and  $|\eta| < 2.5$  that pass the same quality and  $z_0$  criteria as those for the muon ID tracks. The track selection is thus less stringent than the muon selection and the track multiplicity counts the two muons as well.

All selections except the trigger requirements are applied to the MC events. The MC events are assigned a weight based on their probability to pass the trigger requirements. The total probability is calculated by considering each muon in the event and the individual probability of the muon to pass the trigger selection. The MC events are also corrected to account for minor differences between data and MC simulation in the muon reconstruction and identification efficiencies by applying  $p_T$ - and

$\eta$ -dependent scale factors. The tracking efficiency in MC simulation [35] is consistent with data and has been confirmed with additional studies of tracking performance in a dense environment [36]. Thus no corrections are applied to the simulation for tracking performance.

Signal and validation regions are defined after the like-sign dimuon preselection described above and using the  $N_{\text{trk}}$  definition. The signal region is defined as follows:

- Leading-muon  $p_T > 100$  GeV, and
- Track multiplicity  $N_{\text{trk}} \geq 30$ .

The validation regions are defined by inverting one or both of the above requirements. Explicitly, the validation regions are split into two types:

- Leading muon satisfies  $40 < p_T < 100$  GeV without any requirement on  $N_{\text{trk}}$ .
- Leading-muon  $p_T > 100$  GeV and  $N_{\text{trk}} < 24$ .

The validation regions are further split into bins in track multiplicity to test the background estimation techniques described in the next section.

#### V. BACKGROUND ESTIMATION

The backgrounds from Standard Model processes are divided into two categories for ease of estimation: processes where the two muons come from correlated decay chains and processes that produce like-sign dimuons in uncorrelated decay chains. Examples of correlated decay chains are the decays of  $t\bar{t}$  events, where there is a fixed branching ratio to obtain like-sign dimuons. The most likely scenario in  $t\bar{t}$  events is where the leading isolated muon arises from the decay of a  $W$  boson from one of the top quarks, and the subleading muon of the same charge comes either from the semileptonic decay of a  $b$  quark from the other top quark, or from the sequential decay  $b \rightarrow cX \rightarrow \mu X'$  of a  $b$ -quark from the same top quark.

The uncorrelated background estimates arise predominantly from the  $W$ +jets process, where the  $W$  boson decay gives rise to the leading isolated muon and the other muon arises from an in-flight  $\pi/K$  decay, or the semileptonic decay of a  $B$  or  $D$  hadron. Processes such as  $Z$ +jets, and single top in the  $s$ - and  $t$ -channels also give rise to uncorrelated backgrounds when the leading isolated muon arises from the vector boson decay, and one of the jets gives rise to the second muon. The second muon is referred to as a “fake” muon in the subsequent discussion of the uncorrelated background estimate.

##### A. Correlated Background Estimates

The following sources of correlated backgrounds are considered:  $t\bar{t}$  production, diboson production, and single-top production in the  $Wt$  channel. Each of the

three correlated backgrounds is estimated from dedicated MC samples. The background from the  $Wt$  process is small and it has been merged with the  $t\bar{t}$  background in the subsequent discussion and presentation. Other possible sources such as  $t\bar{t}W$  or  $tZ$  production and backgrounds from charge misidentification of muons were found to be negligible.

The sources of uncertainty on the  $t\bar{t}$  background are the choice of MC event generator and parton showering model, the amount of initial- and final-state radiation (ISR/FSR), and the theoretical uncertainty on the production cross section. The  $t\bar{t}$  cross section used is  $\sigma_{t\bar{t}} = 238^{+22}_{-24}$  pb for a top-quark mass of 172.5 GeV. It has been calculated at approximate next-to-next-to-leading order (NNLO) in QCD with HATHOR v1.2 [37] using the MSTW 2008 90% NNLO PDF sets. It incorporates PDF+ $\alpha_S$  uncertainties, according to the MSTW prescription [38], added in quadrature with the scale uncertainty and has been cross-checked with the calculation of Cacciari et al. [39] as implemented in TOP++ v1.0 [40]. The uncertainty on the parton showering model is assessed by comparing the nominal  $t\bar{t}$  prediction with a prediction made using a POWHEG+HERWIG sample. The generator uncertainty is assessed by comparing the POWHEG prediction with predictions made using the ALPGEN and MC@NLO samples. The ISR/FSR uncertainty is determined by using the ACERMC [41] generator interfaced to PYTHIA, and by varying the ISR and FSR scale  $\Lambda_{\text{QCD}}$ , as well as the ISR and FSR cutoff scale. The effect of the top quark mass is studied by generating dedicated samples with top-quark masses of 170 GeV and 175 GeV and is found to be negligible.

The diboson backgrounds have an uncertainty of 6% on the production cross section [42] and a combined generator and parton-showering uncertainty of 24% based on comparisons between SHERPA and POWHEG, and from renormalization and factorization scale variations [43].

In addition to the uncertainties described above, uncertainties from the measurement of trigger efficiency, the muon reconstruction and identification (including uncertainties due to muon  $p_T$  resolution), and the tracking efficiency are considered for each background along with the uncertainty on the integrated luminosity (2.8%). The total systematic uncertainties on the final background estimates from the different sources are summarized in Table I.

## B. Uncorrelated Background Estimates

The uncorrelated background is estimated from data by first measuring the probability for a track to be reconstructed as a muon in a control sample. This probability is then applied to data events with one muon and at least one track to predict the number of dimuon events. This probability is referred to as a fake rate in the subsequent discussion, and the background estimate is referred to as the  $\mu$ +fake background.

TABLE I: The systematic uncertainties on the event yields in the signal region for the different backgrounds and sources, in percent. The uncertainties on signal acceptance are also summarized in the table.

Source	$\mu$ +fake	$t\bar{t}$	Diboson	Signal
Fake rate measurement	34			
Photon trigger	13			
Prompt correction	18			
ISR/FSR		0.7		
Parton showering		9		
Generator		11	24	
Cross section		10	6	
Muon trigger		1.2	1.3	1.3
Muon reconstruction		2.0	1.2	2.3
Luminosity		2.8	2.8	2.8
Tracking efficiency		10	10	10
Fiducial efficiency				15
PDF (Acceptance)				5
Total	41	21	27	19

The fake rate is measured in a control sample consisting of photon+jet events. These events are collected by a single-photon trigger with a threshold at 40 GeV on the photon transverse momentum. The trigger is prescaled, and the collected dataset corresponds to an integrated luminosity of 56  $\text{pb}^{-1}$ . The photon is required to have  $p_T > 45$  GeV, and to satisfy the requirements of Ref. [44]. The photon is also required to satisfy  $E_T^{\Delta R \leq 0.4} < 5$  GeV, where  $E_T^{\Delta R \leq 0.4}$  is the sum of transverse energies of cells in the electromagnetic and hadronic calorimeters in a cone of 0.4 around the photon axis (excluding the cells associated with the photon). The denominator for the fake rate measurement is the number of events with one photon and at least one track. The track must satisfy all the requirements imposed on an ID track associated with a muon as described in Sec. IV. The track is required to be separated from the photon by  $\Delta R > 0.4$ . The numerator is the subset of these events that have at least one muon passing all the criteria associated with the sub-leading muon as described in Sec. IV. In events with more than one track (muon), the track (muon) with the highest  $p_T$  is chosen. The fake rate can have contributions from processes such as  $W(\mu\nu)\gamma$  and  $Z(\mu\mu)\gamma$  that produce prompt muons and bias the fake rate measurement. This prompt-muon bias is corrected by subtracting these contributions based on MC samples generated using SHERPA. The prompt muon correction ranges from approximately 1% at muon  $p_T = 15$  GeV to 30% at  $p_T = 100$  GeV.

The fake rate is parameterized as a function of the  $p_T$  and  $\eta$  of the track, and as a function of  $N_{\text{trk}}$ . Since the signal black hole models produce isolated photon events as well, the fake rate is measured by requiring  $N_{\text{trk}} < 10$  to reduce any potential signal contamination of this control sample. The  $N_{\text{trk}}$  dependence is parame-

TABLE II: The predicted backgrounds in the validation regions compared to the number of observed data events. The uncertainties shown on the total background are the combined statistical and systematic uncertainties.

$N_{\text{trk}}$	$t\bar{t}$	Diboson	$\mu$ +fake	Total	Data
40 GeV < Leading-muon $p_T$ < 100 GeV					
$N_{\text{trk}} < 10$	10000	800	20000	$31000 \pm 4000$	28988
$10 \leq N_{\text{trk}} < 20$	800	3	400	$1200 \pm 100$	1103
$N_{\text{trk}} \geq 20$	16	0.1	6.8	$23 \pm 3$	12
Leading-muon $p_T \geq 100$ GeV					
$N_{\text{trk}} < 10$	2400	140	2300	$4800 \pm 600$	4428
$10 \leq N_{\text{trk}} \leq 11$	190	3	76	$270 \pm 31$	271
$12 \leq N_{\text{trk}} \leq 14$	133	1.1	42	$176 \pm 21$	167
$15 \leq N_{\text{trk}} \leq 19$	60	0.3	17	$77 \pm 9$	68
$20 \leq N_{\text{trk}} \leq 24$	10	0.1	2.9	$13 \pm 2$	13

terized with a linear fit, and extrapolated for all events with  $N_{\text{trk}} > 10$ . The average fake rate based on the criteria defined here is approximately 1%. The fake rate is consistent with that obtained from photon+jet or  $W$ +jet MC samples. The final  $\mu$ +fake background estimate is obtained by selecting events with one muon satisfying the requirements of a leading muon, and one track of the same charge satisfying the requirements associated with an ID track. These  $\mu^\pm \text{track}^\pm$  events are then assigned a weight based on the fake rate calculated for the track, and are then taken through the rest of the analysis chain in the same way as  $\mu^\pm \mu^\pm$  events, with the track acting as the proxy for the subleading muon. There is a correction to the fake estimate from overcounting due to  $t\bar{t}$  and diboson events populating the  $\mu^\pm \text{track}^\pm$  events in data. This correction is estimated from MC simulation to be 2% and is negligible compared to the systematic uncertainty on the fake estimate.

The uncertainties in the  $\mu$ +fake background estimate arise from the statistical uncertainties in measuring the fake rate, the choice of the photon trigger used for the control sample, and the prompt-muon bias correction. The statistical uncertainty in the fake rate is propagated to the final background estimate, along with the uncertainty in the fit parameters used to parameterize the  $N_{\text{trk}}$  dependence. The fake rate is remeasured using data collected by a single-photon trigger with a  $p_T$  threshold of 80 GeV, and the background estimate is recalculated to assess the uncertainty due to the photon trigger. To assess the uncertainty due to the prompt correction, the correction is varied by  $\pm 15\%$  of its nominal value to obtain “up” and “down” fake rates. The final  $\mu$ +fake background estimate is calculated with the up and down fake rates, and the larger variation from the nominal estimate is assigned as a systematic uncertainty. The choice of  $\pm 15\%$  is motivated by the uncertainties described in Ref. [45] that include experimental uncer-

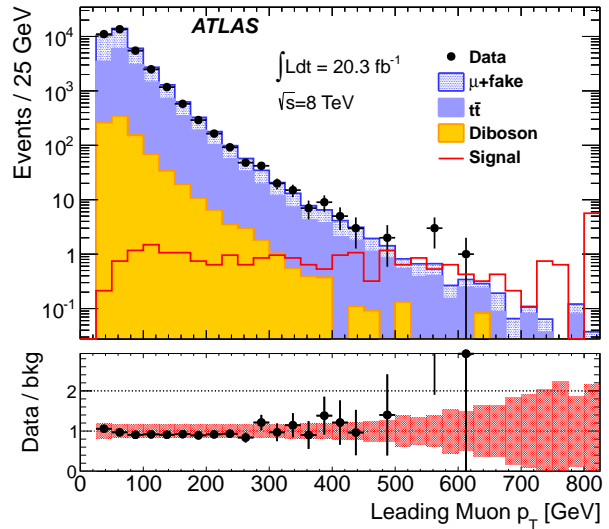


FIG. 1: The leading-muon  $p_T$  distribution for the predicted background and observed data for all like-sign dimuon events passing the preselection criteria. The background histograms are stacked. The signal histogram is overlaid. The last bin along the  $x$ -axis shows the overflows. The bottom panel shows the ratio of data to the expected background (points) and the total uncertainty on the background (shaded area).

tainties on photon reconstruction and identification, and theoretical uncertainties on the production cross sections of the  $W\gamma/Z\gamma$  processes. Table I shows the effect of the systematic uncertainties on the final  $\mu$ +fake background estimate.

## VI. RESULTS AND INTERPRETATION

The background estimation techniques described in the previous section are tested in the validation regions defined in Sec. IV. Table II shows the predicted backgrounds in the various validation regions and the observed yields. The signal contamination is negligible in all the validation regions. Overall, good agreement is observed between the prediction and the observation.

Figure 1 shows the leading-muon  $p_T$  distribution for all like-sign dimuon events (satisfying the preselection). Figure 2 shows the dimuon invariant mass distribution after imposing the  $p_T > 100$  GeV requirement on the leading muon. Figure 3 shows the distribution of the dimuon azimuthal separation,  $\Delta\phi_{\mu\mu}$  for events with  $N_{\text{trk}} \geq 10$  and where the leading muon has  $p_T > 100$  GeV. Figure 4 shows the track multiplicity distribution, which shows good agreement between predicted backgrounds and data. The predicted background and the observed data events in the signal region are shown in Table III. The figures and the table show the expected signal contribution from rotating black holes in a model with  $n = 4$ ,  $M_{\text{TH}} = 5$  TeV and  $M_{\text{D}} = 1.5$  TeV.

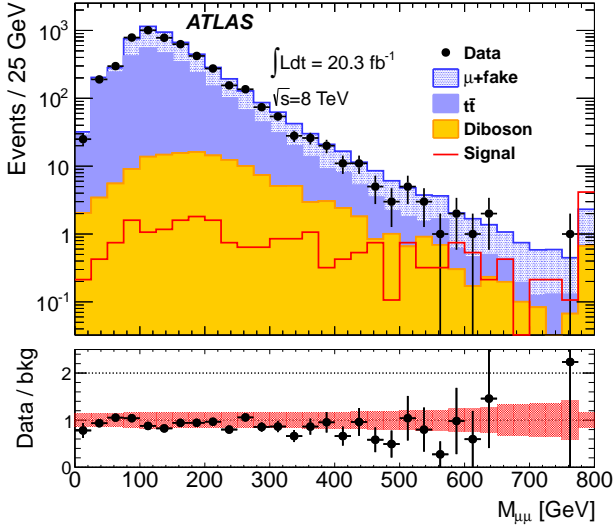


FIG. 2: The dimuon invariant mass ( $M_{\mu\mu}$ ) distribution for the predicted background and observed data for like-sign dimuon events where the leading muon satisfies  $p_T > 100$  GeV. The background histograms are stacked. The signal histogram is overlaid. The last bin along  $x$ -axis shows the overflows. The bottom panel shows the ratio of data to the expected background (points) and the total uncertainty on the background (shaded area).

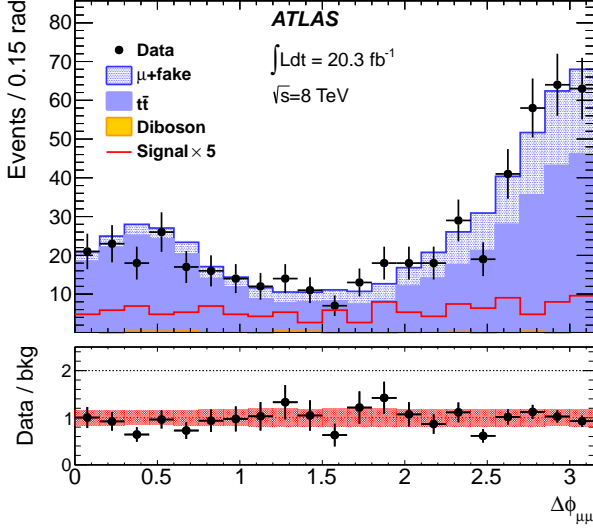


FIG. 3: The dimuon azimuthal separation ( $\Delta\phi_{\mu\mu}$ ) distribution for the predicted background and observed data for like-sign dimuon events where the leading muon satisfies  $p_T > 100$  GeV, and with  $N_{\text{trk}} \geq 10$ . The background histograms are stacked. The signal histogram is overlaid. The bottom panel shows the ratio of data to the expected background (points) and the total uncertainty on the background (shaded area).

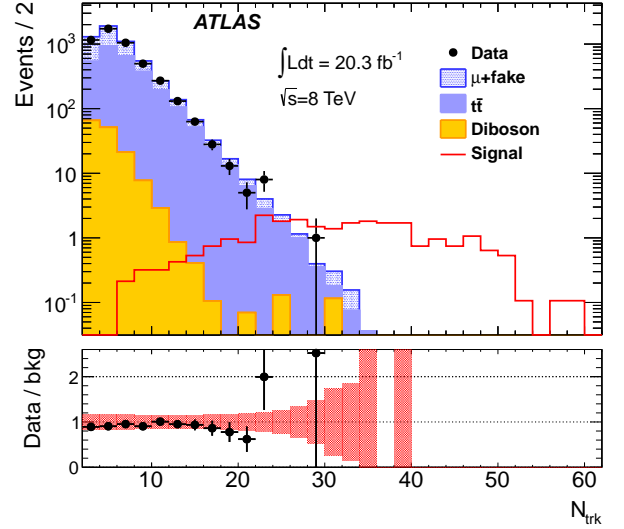


FIG. 4: The track multiplicity ( $N_{\text{trk}}$ ) distribution ( $p_T^{\text{trk}} > 10$  GeV) for the predicted background and observed data for events where the leading muon has  $p_T > 100$  GeV. The background histograms are stacked. The signal histogram is overlaid. The bottom panel shows the ratio of data to the expected background (points) and the total uncertainty on the background (shaded area).

TABLE III: The predicted background in the signal region compared to the number of observed events in data. The MC predictions are shown together with statistical and systematic uncertainties. The expected signal contribution from rotating black holes in a model with  $n = 4$ ,  $M_{\text{TH}} = 5$  TeV and  $M_D = 1.5$  TeV is also shown.

Source	Signal Region
$\mu$ +fake	$0.21 \pm 0.09 \pm 0.09$
$t\bar{t}$	$0.22 \pm 0.08 \pm 0.04$
Diboson	$0.12 \pm 0.08 \pm 0.03$
Total	$0.55 \pm 0.15 \pm 0.10$
Data	0
Signal	$14.2 \pm 1.3 \pm 2.7$

No events are observed in the signal region, which is consistent with the Standard Model prediction. This result is used to set upper limits on the number of events from non-Standard Model sources. The  $CL_s$  method [46] is used to calculate 95% CL upper limits on  $\sigma_{\text{vis}} = \sigma \times BR \times A \times \epsilon$ , where  $\sigma_{\text{vis}}$  is the visible cross section,  $\sigma$  is the total cross section,  $BR$  is the inclusive branching ratio to like-sign dimuons,  $A$  is the acceptance, and  $\epsilon$  is the reconstruction efficiency for non-Standard Model contributions in this final state in the signal region. The observed 95% CL limit on  $\sigma_{\text{vis}}$  is 0.16 fb. The observed limit agrees well with the expected limit of 0.16 fb. The standard deviation ( $\sigma$ ) bands on the expected limit at  $1\sigma$  and  $2\sigma$  are 0.15–0.22 fb and 0.15–0.29 fb, respectively.



Exclusion contours in the plane defined by  $M_{\text{TH}}$  and  $M_{\text{D}}$  for rotating and non-rotating black holes for  $n = 2, 4$ , and  $6$  are obtained. No theoretical uncertainty on the signal prediction is assessed, i.e. the exclusion limits are set for the exact benchmark models as described in Sec. III.

The signal acceptance is measured from the event generator (truth) by imposing the following selections at the particle level. Each event must have at least two true muons with  $p_{\text{T}} > 15$  GeV and  $|\eta| < 2.4$ , and the leading two muons in  $p_{\text{T}}$  must have the same charge. The leading muon must satisfy  $p_{\text{T}} > 100$  GeV. The leading-muon truth-isolation ( $I_{\text{gen}}$ ) is defined as the sum of  $p_{\text{T}}$  of all charged particles with  $p_{\text{T}} > 1$  GeV within a cone of  $\Delta R = 0.2$  around the muon (excluding the muon). The leading muon is required to satisfy  $I_{\text{gen}} < 0.25 \times p_{\text{T}}$ . Each event must also have at least 30 charged particles satisfying  $p_{\text{T}} > 10$  GeV and  $|\eta| < 2.5$ . The ratio of events passing these selections at particle level to the total number of generated events gives the acceptance. The acceptance varies from 11% to 0.2% across the range of model parameters considered here.

The acceptance is then corrected to take into account detector effects. The correction factor,  $\epsilon_{\text{fid}}$ , is defined as the ratio of number of events passing the selection criteria after full detector reconstruction to the number of events passing the acceptance criteria at the particle level. The factor is found to be independent of the number of extra dimensions, and is linearly dependent on  $k = M_{\text{TH}}/M_{\text{D}}$ . The linear dependence is assessed separately for rotating and non-rotating black holes by a fit to the efficiency as a function of  $k$ . For rotating (non-rotating) signals  $\epsilon_{\text{fid}}$  rises from 0.35 (0.3) for  $k = 1$  to 0.55 (0.65) for  $k = 3$ .

The uncertainty on the signal prediction has the following components: the uncertainty on the  $\epsilon_{\text{fid}}$  fit parameters, the uncertainty on luminosity, the uncertainty on acceptance due to the PDFs, the experimental uncertainty on acceptance due to muon trigger and identification efficiencies, and the uncertainty due to tracking efficiency. The uncertainty on acceptance due to PDF was estimated by using the 40 error sets associated to the MSTW 2008 LO PDF set. In the signal region, at high  $N_{\text{trk}}$ , it is possible for small differences between the track reconstruction efficiencies in data and simulation to be magnified. The effect of any possible disagreement between data and simulation is studied by artificially increasing the disagreement and probing the subsequent effect on the signal acceptance. A disagreement of 2% in the per-track reconstruction efficiency translates to a 5% uncertainty in the signal acceptance for  $N_{\text{trk}} \geq 30$ . As a conservative choice, a 10% uncertainty on signal acceptance is assigned to account for possible disagreements in data and simulation track reconstruction efficiency. The uncertainties are summarized in Table I.

Figure 5 shows the expected and observed exclusion contours for non-rotating black holes for  $n = 2, 4$ , and  $6$ . Figure 6 shows the same for rotating black holes. In both figures, the  $1\sigma$  uncertainty band on the expected

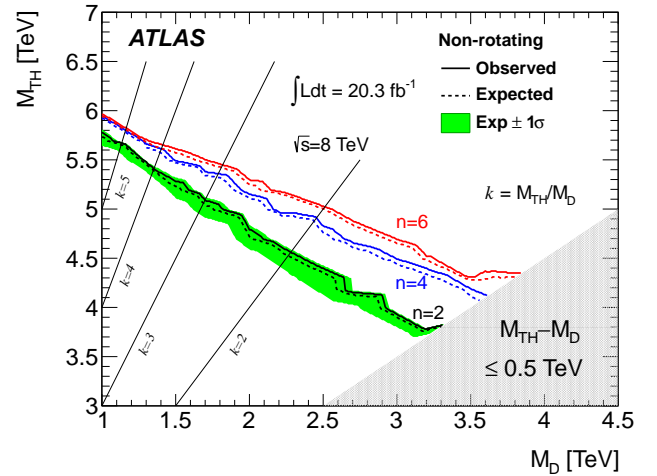


FIG. 5: The 95% CL exclusion contours for non-rotating black holes in models with  $n = 2, 4$ , and  $6$ . The dashed lines show the expected exclusion contour, the solid lines show the observed exclusion contour. The regions below the contours are excluded by this analysis. The  $1\sigma$  uncertainty on the expected limit for  $n = 2$  is shown as a band. Lines of constant slope  $k = M_{\text{TH}}/M_{\text{D}} = 2, 3, 4$ , and  $5$  are also shown. Only slopes  $k \gg 1$  correspond to physical models.

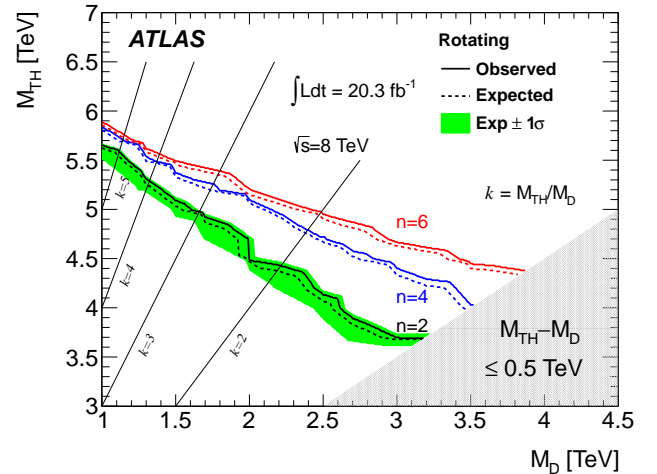


FIG. 6: The 95% CL exclusion contours for rotating black holes in models with  $n = 2, 4$ , and  $6$ . The dashed lines show the expected exclusion contour, the solid lines show the observed exclusion contour. The regions below the contours are excluded by this analysis. The  $1\sigma$  uncertainty on the expected limit for  $n = 2$  is shown as a band. Lines of constant slope  $k = M_{\text{TH}}/M_{\text{D}} = 2, 3, 4$ , and  $5$  are also shown. Only slopes  $k \gg 1$  correspond to physical models.

limit is shown for  $n = 2$ . For each value of  $n$ , the observed limit lies within the  $1\sigma$  band. Lines of constant slope ( $k = M_{\text{TH}}/M_{\text{D}}$ ) of 2, 3, 4 and 5 are also shown. The semi-classical approximations used for black hole production and decay are expected to be valid only for large slopes. The effect of choosing a different set of PDFs



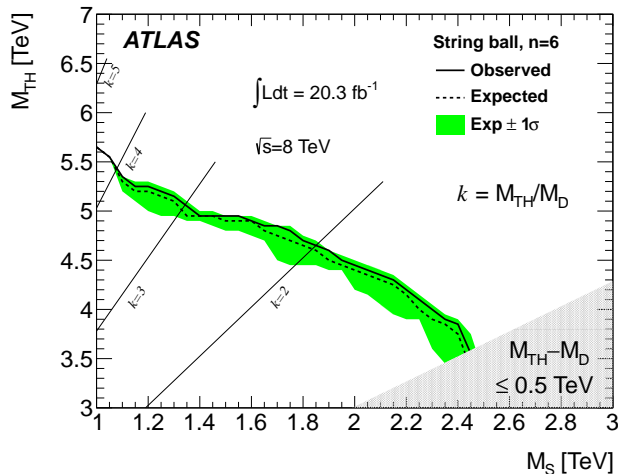


FIG. 7: The 95% CL exclusion contours for string balls in models with  $n = 6$ . The dashed line shows the expected exclusion contour with the  $1\sigma$  uncertainty shown as a band. The solid line shows the observed exclusion contour. The region below the contour is excluded by this analysis. Lines of constant  $k = M_{\text{TH}}/M_{\text{D}} = 2, 3, 4$ , and  $5$  are also shown. Only the values  $k \gg 1$  correspond to physical models.

for signal generation has been studied by considering the CT10 PDF set. The predicted cross sections with the CT10 PDF set are approximately 20% higher, but this has a negligible impact on the exclusion contours due to the rapidly falling cross section with mass for black hole production.

The theory of large extra dimensions can be embedded into weakly-coupled string theory [47, 48], giving rise to string balls whose decay would be experimentally similar to the decay of black holes. Models of string balls have two additional parameters  $M_{\text{S}}$  and  $g_{\text{S}}$ , the string scale and the string coupling constant respectively, in addition to  $M_{\text{TH}}$ ,  $M_{\text{D}}$ , and  $n$ . BLACKMAX is used to simulate the production and decay of string balls, and to obtain exclusion contours in the plane defined by  $M_{\text{TH}}$  and  $M_{\text{S}}$ . Following Ref. [48], the values of  $g_{\text{S}}$  and  $M_{\text{D}}$  are set by  $g_{\text{S}}^2 = 1/5^{\frac{n+2}{n+1}}$ , and  $M_{\text{D}} = 5^{\frac{1}{n+1}} M_{\text{S}}$ . The exclusion contour in the  $M_{\text{TH}}-M_{\text{S}}$  plane for string balls is shown in Fig. 7 for models with  $n = 6$  (where  $g_{\text{S}} = 0.40$  and  $M_{\text{D}} = 1.26 M_{\text{S}}$ ). Table IV shows the summary of lower limits placed on the mass of microscopic black holes and string balls for  $M_{\text{D}} = 1.5$  TeV for different values of  $n$ .

## VII. CONCLUSIONS

A search for microscopic black holes has been carried out using  $20.3 \text{ fb}^{-1}$  of data collected by the ATLAS detector in 8 TeV proton-proton collisions at the LHC. No excess of events over the Standard Model background expectations is observed in the final state with a like-sign dimuon pair and high track multiplicity. Exclusion con-

TABLE IV: The lower limits on  $M_{\text{TH}}$  at 95% CL are summarized for different models.  $M_{\text{D}}$  is fixed at 1.5 TeV. For the string ball model,  $g_{\text{S}} = 0.40$  and  $M_{\text{S}} = M_{\text{D}}/1.26 = 1.2$  TeV.

Model	$n$	$M_{\text{TH}}[\text{TeV}] \geq$
Non-rotating black hole	2	5.3
Non-rotating black hole	4	5.6
Non-rotating black hole	6	5.7
Rotating black hole	2	5.1
Rotating black hole	4	5.4
Rotating black hole	6	5.5
String ball	6	5.3

tours in the plane of the fundamental Planck scale  $M_{\text{D}}$  and the threshold mass  $M_{\text{TH}}$  of black holes are shown and a limit of  $0.16 \text{ fb}$  at 95% CL is set on the visible cross section for any new physics in the signal region defined by a like-sign dimuon pair and high track multiplicity selection.

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angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle  $\theta$  as  $\eta = -\ln \tan(\theta/2)$ .

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